

LETTER TO THE EDITOR

# Detection of a new methanol maser line with ALMA

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Received 9 August 2017 / Accepted 24 September 2017

## ABSTRACT

**Aims.** We investigated the structure and kinematics of the gaseous disk and outflows around the massive young stellar object S255 NIRS3 in the S255IR-SMA1 dense clump.

**Methods.** Observations of the S255IR region were carried out with ALMA at two epochs in the compact and extended configurations.

**Results.** We serendipitously detected a new, never predicted, bright maser line at about 349.1 GHz, which most probably represents the CH<sub>3</sub>OH 14<sub>1</sub>–14<sub>0</sub> A<sup>+</sup> transition. The emission covers most of the 6.7 GHz methanol maser emission area of almost 1'' and shows a velocity gradient in the same sense as the disk rotation. No variability was found in the time interval of several months. The emission is classified as Class II maser and probably originates in a ring at a distance of several hundred AU from the central star.

**Key words.** ISM: molecules – masers – ISM: individual objects: S255IR

## 1. Introduction

Observations of cosmic masers are very important for investigations of physical properties and kinematics of associated astronomical objects as well as for better understanding of the maser excitation mechanisms. Many masing transitions have been discovered so far, but new transitions, especially unpredicted transitions, are always important for the tasks mentioned above. One type of objects where masers are frequently observed and of which masers provide valuable information are star-forming regions, in particular, regions of high-mass star formation. We recently obtained a new data set at 0.8 mm with ALMA toward a bright representative object of this type, S255IR, which is located at a distance of  $1.78^{+0.12}_{-0.11}$  kpc (Burns et al. 2016).

This object has been a target for many investigations. It consists of several dense clumps. The most prominent one, S255IR-SMA1, represents a rotating hot core, probably a disk around the massive ( $\sim 20 M_{\odot}$ ) young stellar object (YSO) S255 NIRS3 (Wang et al. 2011; Zinchenko et al. 2012, 2015). Recently, this YSO attracted greater attention as a result of the discovery of its 6.7 GHz methanol maser flare in 2015 (Fujisawa et al. 2015) and of the subsequent detection of the accretion burst at near-infrared (NIR) wavelengths (Caratti o Garatti et al. 2017; Stecklum et al. 2016).

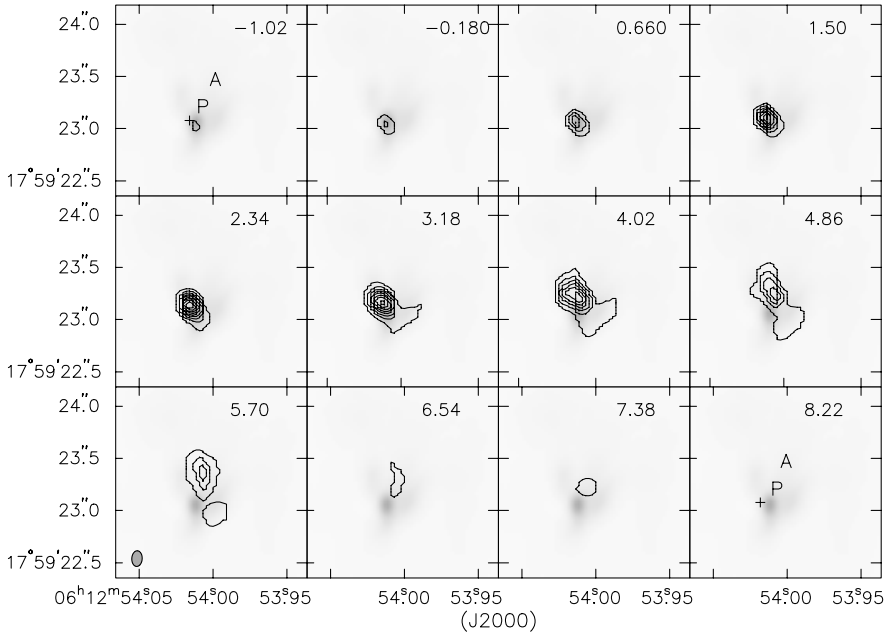
In our ALMA data we noted a very bright apparently non-thermal line toward this object. Here we present the details of the observations and discuss the results.

## 2. Observations

Observations of the S255 IR region were carried out with ALMA in Band 7 at two epochs during Cycle 4 under the project 2015.1.00500.s. The first observing epoch on 2016 April 21, carried out with the 12-m array in its C36-2/3 configuration with baselines ranging between 15 m and 612 m, has 42 antennas

online within a single execution block. The second epoch of observations on 2016 September 9, with the 12-m array in C40-6 configuration and baselines ranging between 15 and 3143 m, contains two execution blocks back-to-back with 39 antennas online. Four frequency division mode (FDM) spectral windows centered at around 335.4 GHz, 337.3 GHz, 349.0 GHz, and 346.6 GHz, with bandwidths of 1875.0 MHz, 234.4 MHz, 937.5 MHz, and 1875.0 MHz, respectively, were set up to cover the continuum and various molecular line features. The spectral feature reported in this paper is situated in the spectral window at 349.0 GHz. A total of 3840 spectral channels are employed to resolve the 937.5 MHz bandwidth, resulting in a spectral channel width of 0.244 MHz and a resolution of 0.488 MHz, or equivalently  $0.42 \text{ km s}^{-1}$  per channel, with Hanning-smoothing applied. For observations taken in the first epoch, J0854+2006 and J0750+1231 were used as the bandpass and flux calibrators, respectively. For the second epoch, J0510+1800 was used for both bandpass and flux calibration purposes. In both epochs, J0613+1708 served as the complex gain calibrator.

All data calibration and reduction was carried out with the Common Astronomy Software Applications (McMullin et al. 2007). Data were first calibrated through the ALMA manual or pipeline calibration process. We further employed self-calibration to improve the imaging quality and dynamical range. A continuum image was first made using visibilities in the wide low-resolution (continuum) spectral window at 335.4 GHz with obvious spectral contamination removed. Two iterations of self-calibration with solution intervals set to 240 s and 60 s were then performed to derive refined phase gains. The final continuum image has a 30% improvement in the dynamic range as compared to the first image. The self-calibration gain solutions were then applied to all spectral windows for further imaging processes. The resulting images achieve an angular resolution of  $0''.10 \times 0''.15$  with Briggs weighting with a robust parameter of 0.5.



**Fig. 1.** Channel map of the CH<sub>3</sub>OH maser emission toward S255IR-SMA1. The CH<sub>3</sub>OH emission is shown in contours, with levels at 0.6, 1.2, 1.8, 2.4, 3.0, 3.6, 4.2, 4.8, and 5.4 Jy/beam. The continuum emission at 335 GHz is displayed in grayscale. The datacube has a channel width of 0.42 km s<sup>-1</sup>, and maps of only every other channel starting from that at -1.02 km s<sup>-1</sup> are shown. Positions of the 5 GHz radio continuum and those of the pre-burst and after-burst CH<sub>3</sub>OH maser clusters reported by Moscadelli et al. (2017) are labeled for reference by the cross, P, and A, respectively, in the first and last channel. The synthesized ALMA beam is shown in the lower left corner.

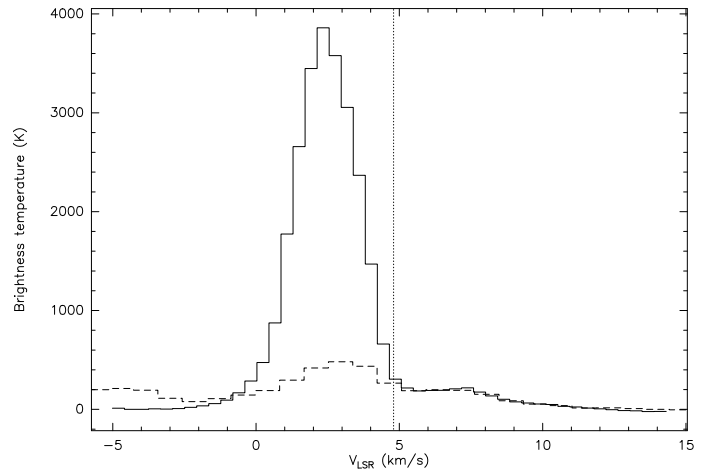
### 3. Results

In the measured spectra toward the SMA1 clump, we see a very strong emission line at about 349.1 GHz. Its peak intensity is higher than the intensity of any other line toward this object by more than an order of magnitude. The brightness reaches about 5.9 Jy/beam, which corresponds to about 3900 K, clearly indicating a nonthermal, most probably masing, nature of this line emission. To the best of our knowledge, a maser emission at this frequency has never been reported before.

An inspection of the major catalogs of spectral lines (Cologne Database for Molecular Spectroscopy, Müller et al. 2001, 2005; and that of the Jet Propulsion Laboratory, Pickett et al. 1998) shows that the most probable identification is the CH<sub>3</sub>OH 14<sub>1</sub>-14<sub>0</sub> A<sup>+</sup> line at 349.1070 GHz with its lower level excitation energy of 243.4483 K. The other nearby lines (within ±5 MHz) belong to complex molecules such as acetone, vinyl cyanide, aminoacetonitrile, and dimethyl ether, which have never been reported to manifest strong maser emission. Methanol has many transitions that are masing at certain conditions. However, the transition mentioned above has never been observed or been predicted to be masing (Sobolev et al. 1997; Cragg et al. 2005; Voronkov et al. 2012).

We display in Fig. 1 the channel map of this emission (by assuming that the rest frequency equals that of the CH<sub>3</sub>OH 14<sub>1</sub>-14<sub>0</sub> A<sup>+</sup> line). The contours in the map run from 0.6 Jy/beam at the lowest level, corresponding to ~400 K, and reach 5.4 Jy/beam, demonstrating the nonthermal nature of the emission. As shown in the figure, strong line emission emerges at about -1 km s<sup>-1</sup> and reaches ~8 km s<sup>-1</sup>, and its peak location gradually moves to the north-north-west.

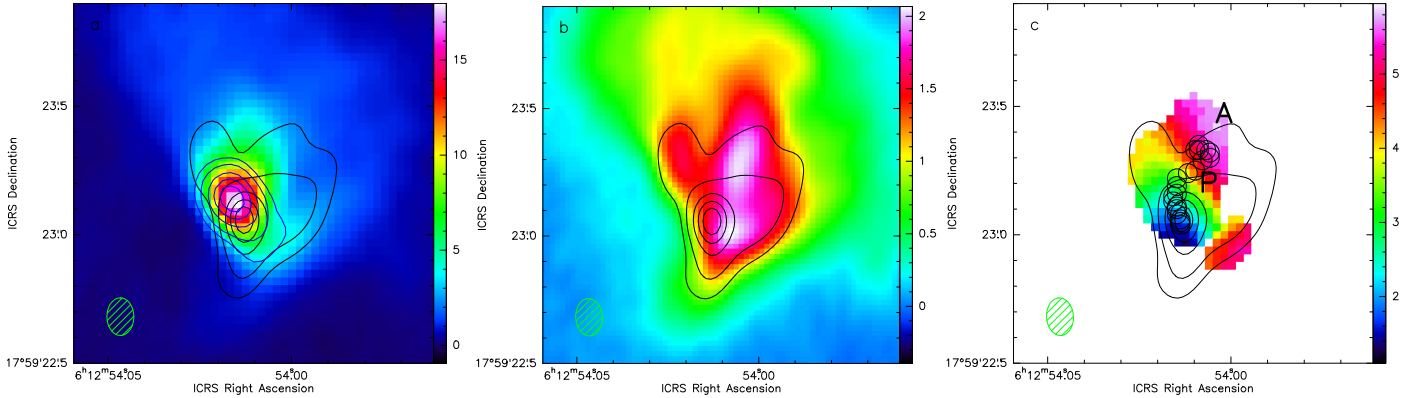
Figure 2 features the spectrum of the maser line toward the brightness peak. For comparison we also show the emission spectrum of one of the apparently thermal methanol lines (the 16<sub>1</sub>-15<sub>2</sub> A<sup>-</sup> line at 345.90397 GHz with the excitation energy of the lower level of 316.0486 K). The profiles of the maser and thermal lines are in general similar, but the blueshifted part of the maser line is strongly amplified. The width of the brightest maser feature is about 2.5 km s<sup>-1</sup>.



**Fig. 2.** Spectrum of the brightness temperature in the masing methanol 14<sub>1</sub>-14<sub>0</sub> A<sup>+</sup> line on the velocity axis (zero velocity corresponds to the rest frequency of this transition) within the synthesized beam toward the maser brightness peak (solid line). The dashed line shows the spectrum centered on the CH<sub>3</sub>OH 16<sub>1</sub>-15<sub>2</sub> A<sup>-</sup> line. The intensity of this line is multiplied by a factor of 3. The dotted vertical line indicates the systemic velocity of the S255IR-SMA1 clump (4.8 km s<sup>-1</sup>), as found by Zinchenko et al. (2015).

The spatial distribution of the emission in this line is very different from that in the thermal methanol lines, as can be seen in Figs. 3a and b, where we also present the map of the methanol 16<sub>1</sub>-15<sub>2</sub> A<sup>-</sup> line for comparison. This line in general follows the continuum emission, while the former appears relatively centrally peaked with its peak slightly shifted NE from the continuum peak. The peak flux density integrated over a circle of about 1'' is about 25 Jy (at a velocity of ~4 km s<sup>-1</sup>).

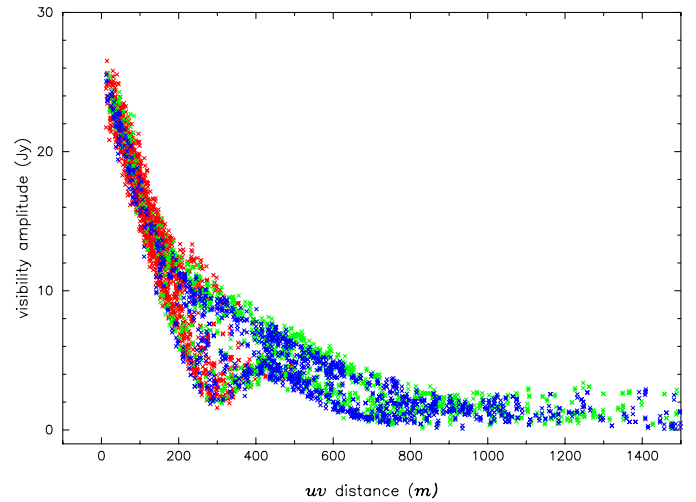
Moscadelli et al. (2017) reported several clusters of the 6.7 GHz maser spots observed with the Jansky Very Large Array and the European VLBI Network. In Fig. 1 the position of their 5 GHz radio continuum as well as the positions of the dominating pre-/post-burst 6.7 GHz methanol maser clusters labeled P and A by Moscadelli et al. (2017) are marked. The emission detected in our observations covers most of their



**Fig. 3.** Maps of the integrated emission in the  $\text{CH}_3\text{OH}$   $14_1-14_0$   $\text{A}^+$  line (a) and  $16_1-15_2$   $\text{A}^-$  line (b). The thick blue contours in the panel a show the map at the velocity of the peak brightness. c) The first-moment map of the maser emission with the intensity threshold of  $0.9 \text{ Jy beam}^{-1}$  (which corresponds to approximately 600 K). The chain of circles shows the positions of the emission peaks in different spectral channels. The locations of the pre-burst and after-burst 6.7 GHz maser clusters are indicated. The thin contours in all panels show the 0.8 mm continuum emission. The contour levels are 0.1, 0.2, 0.4, 0.6, and 0.8 of the peak intensity. The synthesized ALMA beam is shown in the lower left corner.

6.7 GHz methanol emission area of almost  $1''$ . The main peak of our emission, however, is closer to one of their (unlabeled) clusters of maser spots northeast of the 5 GHz continuum. There is a secondary emission peak, southwest of the main peak, which coincides with another unlabeled cluster of the maser spots in Moscadelli et al. (2017). The emission in individual spectral channels is much more compact. At the velocity of the peak brightness ( $\sim 2.34 \text{ km s}^{-1}$ ), it is about  $0''.17 \times 0''.08$  ( $\sim 300 \times 140 \text{ AU}$ ). There is a significant velocity gradient across the emission region, as noted earlier. This gradient, best demonstrated by the first-moment map of the emission in Fig. 3c, is in the same sense as the rotation of the disk (Wang et al. 2011; Zinchenko et al. 2015).

Maser emission is known to be time variable, and we examined this possibility given that our two epochs of observations are separated by about five months. Nevertheless, the array configurations and hence the angular resolutions in the two observational epochs are different; a direct comparison of the observed flux densities in the image domain through separate imaging of the different individual observational datasets is not straightforward. We instead compare the visibility amplitudes of the maser emission observed at different epochs. Figure 4 displays the visibility amplitudes of the maser emission against the baseline lengths. The visibilities are taken from the integrated emission peak channel at around  $4 \text{ km s}^{-1}$  from the three execution blocks (one from 2016 April 21 in red, and two from 2016 September 9 in green and blue). As shown in the figure, the visibilities from the three datasets have a fully consistent amplitude distribution within baseline lengths shorter than about 180 meters, which corresponds to  $205 \text{ k}\lambda$ , or an angular scale of  $\sim 1''$ . This implies no noticeable flux density variation at an angular scale  $\geq 1''$ . The visibility amplitudes at intermediate baseline lengths (between 180 m and  $\sim 600$  m) appear to primarily have two branches of amplitudes, with those from 2016 April mainly in the higher branch and those from September in the lower branch. We verified that the reason is that these visibilities from the two epochs sample different parts (direction) of the  $uv$  domain. Still, the amplitude scales of the overlapping region are fully consistent. The consistency of amplitudes between visibilities from different epochs also preserve the methanol emission channels throughout, suggesting no line profile variation between the observations. We have no sampling at long baselines in April to differentiate or identify any flux density variation at angular scales significantly shorter than  $0''.4$ . Given that



**Fig. 4.** Visibility amplitudes of the maser emission against the baseline lengths. The visibilities are taken from the maser emission peak channel at  $4 \text{ km s}^{-1}$  from the three execution blocks (one from 2016 April 21 in red, and two from 2016 September 9 in green and blue). Only visibilities with baseline lengths shorter than 1500 m are shown.

the absolute flux calibration for band 7 is at a level of about 5% based on the ALMA Cycle 4 Technical Handbook, we conclude that the methanol emission shows no variation above this level in its total flux density over a timescale of several months.

## 4. Discussion

Analysis of observational data has shown that there are two classes of methanol masers, Class I and Class II, which emit in different transition sets (Menten 1991). Masers of one class never display strong emission in transitions characteristic of the other class. Observational results are supported by theoretical modeling that shows that the warm radiation field cools down collisionally excited maser transitions and provides power to transitions that are suppressed by the pumping mechanism with collisional excitation (Sobolev et al. 2007; Voronkov et al. 2005; Cragg et al. 1992). As a result, Class II methanol masers reside closer to YSOs and trace circumstellar disks and inner parts of the outflows, while Class I masers are associated with more distant parts of the outflows and shocked regions. This conclusion

finds observational support for example in [Sanna et al. \(2015, 2017\)](#) and [Voronkov et al. \(2014\)](#).

Although methanol masers of different classes are formed under different conditions, they can be located close to each other or even coincide due to projection effects (e.g., [Kurtz et al. 2004](#); [Ellingsen 2005](#); [Sobolev et al. 2007](#)). This means that knowledge about the class of maser transition is important for understanding which parts of the YSO vicinity are traced.

Existence of two classes of methanol masers can be qualitatively explained by selection rules for radiative transitions, which lead to overpopulation of the sets of energy levels with particular value of quantum number  $K$  (called  $K$ -ladders) with respect to the levels from the different  $K$ -ladders ([Menten 1991](#)). As a result, maser transitions form series of transitions that differ by the value of the quantum number  $J$ . The 349.1070 GHz transition under study belongs to the  $J_1-J_0$   $A^{++}$  series. Maser emission in transitions of this series was not observed or theoretically predicted. However, qualitative considerations show that this transition probably belongs to Class II. Observations show that the strongest Class II maser occurs in the  $5_1-6_0$   $A^+$  transition, while the strongest Class I maser occurs in the  $7_0-6_1$   $A^+$  transition. Class II methanol masers are therefore characterized by overpopulation of  $A^+$  levels of the  $K = 1$  ladder with respect to the  $A^+$  levels of the  $K = 0$  ladder, while for Class I masers the situation is opposite. Model calculations for Class II methanol maser conditions by [Cragg et al. \(2005\)](#) show that  $A^-$  levels of the  $K = 1$  ladder tend to be overpopulated with respect to  $A^+$  levels of the  $K = 1$  ladder. This means that the 349.1070 GHz transition from  $J_1-J_0$   $A^{++}$  series most probably belongs to the Class II family. Strong support for this conclusion comes from the observational evidence that emission in the newly discovered maser transition covers the same velocity range and shows a similar position-velocity dependence as the strongest Class II methanol maser at 6.7 GHz ([Moscadelli et al. 2017](#)). The spectra and the peak locations of the new maser may not coincide with those of the 6.7 GHz masers because of the different sensitivity to the physical conditions, which is usual for Class II methanol maser transitions ([Cragg et al. 2005](#)).

The  $J_1-J_0$   $A^{++}$  line series contains transitions with frequencies starting from 303.37 GHz. Many of them can be observed with the same instrument, which allows avoiding calibration problems in determining the relative line intensities. The lower levels of these transitions cover a rather wide range of energies from 2.32 K for the  $1_0$   $A^+$  level to 243.45 K for the  $14_0$   $A^+$  lower level of the transition under study and higher. It is worth mentioning that in our SMA data obtained in December 2010, the  $\text{CH}_3\text{OH}$   $13_1-13_0$   $A^{++}$  line at 342.729796 GHz was detected in quasi-thermal emission ([Zinchenko et al. 2015](#)). This may indicate that the discovered maser is related to the accretion burst in 2015, but additional modeling and monitoring of these lines is needed for an exact answer. Observations of the methanol line series proved to be a good tool for the physical parameter diagnostics both in maser and quasi-thermal cases ([Voronkov et al. 2006](#); [Menten et al. 1988](#); [Salii & Sobolev 2006](#)). This is very important because the new maser line was detected toward positions that are spread over the disk.

An inspection of our channel maps shows that the peaks of the maser emission at different velocities lie on a curve that resembles an arc (Fig. 3c). This can be part of a ring in the disk or on the surface of the disk at a substantial distance of several hundred AU from the star. Taking into account the disk orientation, which follows from the outflow data ([Zinchenko et al. 2015](#); [Burns et al. 2016](#)), it becomes clear that the brightest emission comes from the part of the disk that is closer to us. For the lower

brightness component, the velocity gradient is less clear, but its location shows that it may originate on the far side of the disk. The difference in brightness may be related to the different background level.

## 5. Conclusions

We reported the discovery of a new, never predicted, rather bright maser line at about 349.1 GHz. Most probably, this line represents the  $\text{CH}_3\text{OH}$   $14_1-14_0$   $A^{++}$  transition with the excitation energy of the lower level of 243.4483 K. It can be classified as a Class II methanol maser.

This line was detected toward the S255IR-SMA1 clump, which represents a rotating disk around the massive YSO S255 NIRS3. The maser emission is observed in the same area as the 6.7 GHz methanol maser and shows a velocity gradient in the same sense as the disk rotation. It probably originates in a ring at a distance of several hundred AU from the central star. No variability was found in the time interval of several months.

*Acknowledgements.* We are grateful to the anonymous referee for the helpful comments. I.Z. was supported by the Russian Science Foundation (project No. 17-12-01256). This paper makes use of the following ALMA data: ADS/JAO.ALMA#2015.1.00500.S. ALMA is a partnership of ESO (representing its member states), NSF (USA) and NINS (Japan), together with NRC (Canada), MOST and ASIAA (Taiwan), and KASI (Republic of Korea), in cooperation with the Republic of Chile. The Joint ALMA Observatory is operated by ESO, AUI/NRAO and NAOJ.

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